Lectures on relativistic gravity and cosmology.

Lectures 1-2

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Special Theory of Relativity

Special relativity kinematics

Special relativity dynamics

Special relativity electrodynamics

The principle of relativity and the maximal velocity

Inertial reference frames (IRF): they exist and move evenly and rectilinearly with respect to each other.

The principle of relativity

All laws of nature are the same in all inertial reference frames.

Hypothesis: there exists a maximal velocity. According to the principle of relativity, it should be the same in all inertial systems. A natural candidate: the light velocity in vacuum:

$$c = 2.99792458 \cdot 10^{10} \, \mathrm{cm/c}$$

The principle of relativity + the maximal velocity = the Special Theory of Relativity (SR).



Classification of areas in physics

Three independent fundamental dimensional physical constants: c, \hbar , G. All other physical constants can be made dimensionless.

- c relativistic physics.
- \hbar quantum physics.
- G gravitational physics.
- c, \hbar relativistic quantum physics.
- c, G relativistic gravitational physics.
- c, \hbar , G relativistic quantum gravitational physics.

Interval

The space-time interval s_{12} between two events:

$$s_{12}^2 = c^2(t_2 - t_1)^2 - (x_2 - x_1)^2 - (y_2 - y_1)^2 - (z_2 - z_1)^2.$$

Its differential form:

$$ds^2 = c^2 dt^2 - dx^2 - dy^2 - dz^2 = g_{ik} dx^i dx^k$$
, $i, k = 0, 1, 2, 3, x^0 = ct$

 g_{ik} - the metric tensor, $g_{il}g^{kl} = \delta_i^k$.

The first conventional choice of sign: space-time signature (+---).

The Einstein rule: $a_i b^i \equiv g_{ik} a^i b^k \equiv \sum_{i=0}^3 a_i b^i$.

For light propagation: $s_{12} = ds^2 = 0$ in all IFR.



For non-zero intervals: $ds^2 = a(|\mathbf{V}|)ds'^2$ due to homogeneity and isotropy of space and homogeneity of time. Let K, K_1 , K_2 - three IRF, and $\mathbf{V_1}$ and $\mathbf{V_2}$ - relative velocities of K_1 and K_2 with respect to K. Then

$$ds^2 = a(V_1)ds_1^2 = a(V_2)ds_2^2$$

On the other hand,

$$ds_1^2 = a(V_{12})ds_2^2$$

where V_{12} is the modulus of velocity of K_2 with respect to K_1 . Thus,

$$\frac{\mathit{a}(V_1)}{\mathit{a}(V_2)} = \mathit{a}(V_{12}) o \mathit{a}(V) \equiv \mathit{const} = 1$$

$$ds^2 = ds'^2$$
, $s = s'$



Light cone

 $s_{12}^2>0$ - a time-like interval, inside the light cone of the event 1 (future light cone for $t_2>t_1$, past light cone for $t_2< t_1$). It is possible to find IRF where both events 1 and 2 occur in the same place. In this IRF, the time difference between these events (called the proper time difference) is $\tau=s_{12}/c$.

 $s_{12}^2 < 0$ - a space-like interval, outside the light cone of the event 1. No absolute simultaneity. Depending on the choice of an IRF, the event 2 can be in the past, in the future or simultaneous to the event 1.

The Lorentz transformation

The Lorentz transformation:

$$x = \gamma(V)(x' + Vt'), \ y = y', \ z = z', \ t = \gamma(V)\left(t' + \frac{V}{c^2}x'\right)$$

where the Lorentz factor is

$$\gamma(V) = \left(1 - \frac{V^2}{c^2}\right)^{-1/2}$$

The Lorentz group O(1,3).

Compare with the Galilean transformation:

$$x = x' + Vt', y = y', z = z', t = t'$$



4-velocity and 4-acceleration

4-velocity:

$$u^i = rac{dx^i}{ds}, \quad u_i u^i = 1, \quad u^i = (\gamma(v), \ \gamma(v) \mathbf{v})$$

$$ds = cd\tau = c\gamma^{-1}(v)dt$$

3-velocity transformation (IRF K' is moving with the velocity V relative to IRF K along the x-axis):

$$v_{x} = \frac{v'_{x} + V}{1 + \frac{v'_{x}V}{c^{2}}}, \quad v_{y} = \frac{v'_{y}}{\gamma(V)\left(1 + \frac{v'_{x}V}{c^{2}}\right)}, \quad v_{z} = \frac{v'_{z}}{\gamma(V)\left(1 + \frac{v'_{x}V}{c^{2}}\right)}$$

where
$$\mathbf{v} = \frac{d\mathbf{r}}{dt}$$
, $\mathbf{v}' = \frac{d\mathbf{r}'}{dt}$.

4-acceleration:

$$w^{i} = \frac{du^{i}}{ds} = \frac{d^{2}x^{i}}{ds^{2}}, \quad w^{i}u_{i} = 0$$



Action

Action for a free massive particle:

$$S = -mc \int_{a}^{b} ds$$

The coefficient is chosen in such a way that in the non-relativistic limit $v \ll c$, this expression would give the correct expression for the Lagrangian equal to the particle kinetic energy (up to a constant):

$$S = \int_{t_1}^{t_2} L dt = -mc \int_{t_1}^{t_2} \frac{dt}{\gamma(v)}, \ L \approx = -mc^2 + \frac{mv^2}{2}$$

Equations of motion

The principle of least action (more accurately, the principle of stationary action): $\delta S = 0$.

Taking into account that $ds = \sqrt{dx_i dx^i}$, we get:

$$\delta S = -mc \int_{a}^{b} \frac{dx_{i} \delta dx^{i}}{ds} = -mc \int_{a}^{b} u_{i} d\delta x^{i} =$$

$$-mcu_{i} \delta x^{i} \Big|_{a}^{b} + mc \int_{a}^{b} \delta x^{i} \frac{du_{i}}{ds} ds$$

Equations of motion follow from the variation of action with respect to arbitrary virtual trajectories with the fixed end-points $(\delta x^i)_a = (\delta x^i)_b = 0$:

$$\frac{du^i}{ds} = 0$$

4-momentum

4-momentum is defined as the coordinate derivative of the action estimated on real trajectories with the fixed initial point only $(\delta x^i)_a = 0$:

$$p_{i} = -\frac{\partial S}{\partial x^{i}} = mcu_{i} = \left(\frac{\mathcal{E}}{c}, -\mathbf{p}\right)$$
$$\mathcal{E} = m\gamma(v)c^{2}, \ \mathbf{p} = m\gamma(v)\mathbf{v} = \frac{\mathcal{E}\mathbf{v}}{c^{2}}$$

$$p^{i}p_{i} = m^{2}c^{2} \rightarrow \mathcal{E}^{2} = p^{2}c^{2} + m^{2}c^{4}$$

In the non-relativistic limit: $\mathcal{E} = mc^2 + \frac{mv^2}{2}$. For the v = 0, the correct form of the famous Einstein formula follows:

$$\mathcal{E} = mc^2$$

The massless limit: photons, gravitons.

$$\mathcal{E} = |\mathbf{p}|c = \hbar\omega = h\nu$$



Relativistic Hamilton-Jacobi equation

$$\frac{\partial S}{\partial x_i} \frac{\partial S}{\partial x^i} = g^{ik} \frac{\partial S}{\partial x^i} \frac{\partial S}{\partial x^k} = m^2 c^2$$

Transition to the non-relativistic Hamilton-Jacobi equation: introduce the new action $S' = S + mc^2t$. For $c \to \infty$:

$$\frac{\partial S'}{\partial t} + \frac{1}{2m} \left[\left(\frac{\partial S'}{\partial x} \right)^2 + \left(\frac{\partial S'}{\partial y} \right)^2 + \left(\frac{\partial S'}{\partial z} \right)^2 \right] = 0$$

The massless limit: light in the geometric optics approximation. The eikonal equation:

$$\frac{\partial \psi}{\partial x_i} \frac{\partial \psi}{\partial x^i} = 0, \quad k_i = \frac{p_i}{\hbar} = -\frac{\partial \psi}{\partial x^i}$$

Particle interactions

$$(\sum_n \mathcal{E}_n)^2 - (\sum_n \mathbf{p}_n)^2 c^2 = \text{inv}$$

A proton with an energy $\mathcal{E}\gg m_pc^2$ is moving through thermal radiation with the characteristic photon energy $\varepsilon_\gamma\sim kT_\gamma\ll\mathcal{E}$. For which energy the process $p+\gamma\to p+\pi_0$ $(n+\pi_+)$ becomes possible $(T_\gamma=2.725~{\rm K})$?

$$(\mathcal{E}+arepsilon_{\gamma})^2-(pc-arepsilon_{\gamma})^2=(m_p+m_\pi)^2c^4 \ \ \mathcal{E}=rac{2m_pm_\pi+m_\pi^2}{4arepsilon_{\gamma}}\,c^4\sim 10^{20}\mathrm{eV}$$

The Greisen-Zatsepin-Kuzmin effect (1966) resulting in the cutoff (GZK cutoff) in the spectrum of high-energy cosmic rays - observed (2007)!

Testing the Lorentz invariance

- 1. Strong interaction: particles with up to $\sim 3 \cdot 10^{20}$ eV in high-energy cosmic rays.
- 2. Electromagnetic and weak interactions.

arXiv:1807.06504 - photons and neutrino events from the blazar TXS 0506+056 ($\varepsilon_{\nu}\approx$ 290 TeV).

Photons: $\mathcal{E}_3 \gtrsim 10^{20}$ GeV, $\mathcal{E}_4 \gtrsim 10^{11}$ GeV (corrections $\frac{p^3c^3}{\mathcal{E}_3}$ and $\frac{p^4c^4}{\mathcal{E}_4^2}$ in the dispersion relation for $\mathcal{E}^2(p)$).

Neutrinos: $\mathcal{E}_3 \gtrsim 10^{19}$ GeV, $\mathcal{E}_4 \gtrsim 10^{11}$ GeV.

3. Gravitational interaction.

GW170817 + GRB 170817A event. Distance: 40 \pm 10 Mpc.

X-ray signal: 1.7 s after the peak of the GR one.

$$\left| \frac{v_{GW}}{c} - 1 \right| < 3 \cdot 10^{-15}$$



Electromagnetic interactions

New physical quantity: electric charge e. Quantized, $\pm n$ for free particles (protons, electrons, etc.), $\pm \frac{n}{3}$ for quarks.

The fine structure constant $\alpha = \frac{e^2}{\hbar c}$. The 2018 CODATA recommended value $\alpha^{-1} = 137.035999084(21)$.

No change in the Maxwell equations: they are already invariant under the Lorentz transformation.

The 4-vector potential: $\mathbf{A}^i = (\phi, \mathbf{A})$.

The electromagnetic field tensor: $\vec{F}_{ik} = A_{k,i} - A_{i,k}$. In the 3D form:

$$\mathbf{E} = -\frac{1}{c} \frac{\partial \mathbf{A}}{\partial t} - \operatorname{grad} \phi, \ \mathbf{H} = \operatorname{rot} \mathbf{A}$$

$$E_x = F_{01}, \ E_y = F_{02}, \ E_z = F_{03}, \ H_x = F_{32}, \ H_y = F_{13}, \ H_z = F_{21}$$

Gauge transformations: $A_k \to A_k' = A_k - \frac{\partial f}{\partial x^k}$

Action and equations for a charge

Action for a point massive charge:

$$S = -mc \int_{a}^{b} ds - \frac{e}{c} \int_{a}^{b} A_{i} dx^{i}$$

Equations of motion:

$$mc\frac{du^{i}}{ds} = \frac{e}{c}F^{ik}u_{k}$$

In the 3D form:

$$\frac{d\mathbf{p}}{dt} = e\mathbf{E} + \frac{e}{c}[\mathbf{vH}], \quad \frac{d\mathcal{E}}{dt} = e\mathbf{E}\mathbf{v}$$

where \mathcal{E} and \mathbf{p} are the same as in the absence of electromagnetic field.

Action and equations for the field

Action or the electromagnetic field:

$$S_{em}=-rac{1}{16\pi c}\int F_{ik}F^{ik}d\Omega,~~d\Omega=c~dt~dx~dy~dz$$

The Maxwell equations (comma - partial derivative):

$$F_{ik,l} + F_{li,k} + F_{kl,i} = 0$$

$$F^{ik}_{,k} = \frac{4\pi}{c} j^i$$

The electric 4-current $j^i = (c\rho, \mathbf{j})$.

Testing the constancy of α .

1. Laboratory measurements (2008):

$$\frac{\dot{\alpha}}{\alpha} = (-1.6 \pm 2.3) \times 10^{-17}$$
 per year.

2. Measurements using remote quasars: $\left|\frac{\Delta\alpha}{\alpha}\right|\lesssim 10^{-5}$ for the last 10-12 billion years.